# 5. Spectral analysis

- Power spectral density
- Periodogram analysis
- Window functions

# **Power Spectral density**

#### Wiener-Khinchin theorem:

If a process is wide-sense stationary, the autocorrelation function and the power spectral density form a Fourier transform pair:

#### **Continuous**

$$S(\omega) = \int_{-\infty}^{\infty} e^{-i\omega\tau} R(\tau) d\tau \quad \iff \quad R(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} S(\omega) e^{i\omega t} d\omega$$

#### **Discrete**

$$S(\omega) = \sum_{k=-\infty}^{k=\infty} R(k)e^{-i\omega k} \quad \Longleftrightarrow \quad R(k) = \frac{1}{2\pi} \int_{-\pi}^{\pi} S(\omega)e^{i\omega k} d\omega$$

(Under a condition for the existence of the Fourier transform, e.g., R(t) is absolutely integrable or R(k) is absolutely summable)

# **Properties of PSD**

- $S(\omega)$  is self-adjoint, i.e.,  $S(\omega) = S^*(\omega), \forall \omega$
- $S(\omega) \succeq 0$  for all  $\omega$

• 
$$\int_{-\infty}^{\infty} S(\omega)d\omega = R(0) = \mathbf{E} x(t)x(t)^* \succeq 0$$
 (average power)

- ullet For real processes,  $S(-\omega)=S(\omega)^T$
- ullet For discrete-time processes,  $S(\omega)$  is a periodic function of period  $2\pi$

### **Cross-power spectral density**

The cross-power spectrum of x(t) and y(t) is the Fourier transform of the cross correlation  $R_{xy}(\tau)$ :

#### **Continuous**

$$S_{xy}(\omega) = \int_{-\infty}^{\infty} e^{-i\omega\tau} R_{xy}(\tau) d\tau \quad \Longleftrightarrow \quad R_{xy}(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} S_{xy}(\omega) e^{i\omega t} d\omega$$

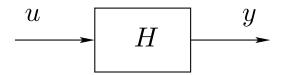
#### **Discrete**

$$S_{xy}(\omega) = \sum_{k=-\infty}^{k=\infty} R_{xy}(k)e^{-i\omega k} \quad \Longleftrightarrow \quad R_{xy}(k) = \frac{1}{2\pi} \int_{-\pi}^{\pi} S_{xy}(\omega)e^{i\omega k}d\omega$$

It follows from  $R_{xy}(-\tau)=R_{yx}^*(\tau)$  that

$$S_{xy}(\omega) = S_{yx}^*(\omega)$$

### LTI systems with random inputs



If u(t) is wide-sense stationary, y(t) is also wide-sense stationary

$$\mathbf{E} y(t) = \sum_{s=-\infty}^{\infty} h(s) \mathbf{E} u(t-s) = \mu_u \sum_{s=-\infty}^{\infty} h(s)$$

The mean is constant for all t

$$R_{y}(t_{1}, t_{2}) = \sum_{s=-\infty}^{\infty} \sum_{v=-\infty}^{\infty} h(s) \mathbf{E}[u(t_{1} - s)u(t_{2} - v)^{*}]h^{*}(v)$$
$$= \sum_{s=-\infty}^{\infty} \sum_{v=-\infty}^{\infty} h(s)R_{u}(t_{1} - t_{2} + v - s)h^{*}(v)$$

 $R_y(t_1,t_2)$  depends only on the time shift  $t_1-t_2$ 

### LTI systems with random inputs

The input-output cross correlation is

$$R_{yu}(t_1, t_2) = \mathbf{E} \sum_{k=-\infty}^{\infty} h(k)u(t_1 - k)u(t_2)^*$$
$$= \sum_{k=-\infty}^{\infty} h(k)R_u(t_1 - t_2 - k)$$

Thus y(t), u(t) are jointly wide-sense stationary with

$$R_{yu}(\tau) = \sum_{k=-\infty}^{\infty} h(k)R_u(\tau - k)$$

It also follows that

$$R_y(\tau) = \sum_{k=-\infty}^{\infty} h(k) R_{uy}(\tau - k)$$

# **Spectral relations for LTI systems**

Using the convolution property of the Fourier transform of  $R_{yu}(\tau), R_y(\tau)$ , we have the relations:

$$S_{yu}(\omega) = H(\omega)S_u(\omega), \quad S_y(\omega) = H(\omega)S_{uy}(\omega)$$

With  $S_{uy}(\omega) = S_{yu}^*(\omega)$ , we have

$$S_y(\omega) = H(\omega)S_u(\omega)H(\omega)^*$$

In terms of z-transform, this could be written as

$$S_y(z) = H(z)S_u(z)H(z)^*$$

where  $H(z)^* = H(\bar{z})^T$  and we should be aware that  $z = e^{\mathrm{i}\omega}$  in the analysis

### Example 1

Suppose the covariance function of a staionary process is given by

$$R(k) = a^{|k|}, \quad |a| < 1$$

The spectral density can be obtained via z-transform

$$S(z) = \sum_{k=-\infty}^{\infty} a^{|k|} z^{-k} = \sum_{k=-\infty}^{-1} a^{-k} z^{-k} + \sum_{k=0}^{\infty} a^{k} z^{-k}$$
$$= \frac{az}{1 - az} + \frac{z}{z - a} = \frac{1 - a^{2}}{(1 - az)(1 - az^{-1})}$$

Substituting  $z = e^{i\omega}$  gives

$$S(\omega) = \frac{1 - a^2}{(1 - ae^{i\omega})(1 - ae^{-i\omega})} = \frac{1 - a^2}{1 + a^2 - 2a\cos\omega}$$

### Example 2

A recursion equation

$$y(t) = ay(t-1) + e(t)$$

where e(t) is a white noise with variance  $\lambda^2$ 

The transfer function is given by

$$H(z) = \frac{1}{1 - az^{-1}}$$

The spectral density of y is therefore

$$S_y(\omega) = \frac{\lambda^2}{(1 - ae^{-i\omega})(1 - ae^{i\omega})} = \frac{\lambda^2}{1 + a^2 - 2a\cos\omega}$$

# **Spectral analysis**

Use the same model as in correlation analysis:

$$R_{yu}(\tau) = \sum_{k=0}^{\infty} h(k)R_u(\tau - k)$$

Taking DFT gives the spectral representation

$$S_{yu}(\omega) = H(\omega)S_u(\omega)$$

If  $S_u(\omega) \succ 0$  for all  $\omega$ , then we can estimate

$$\hat{H}(\omega) = \hat{S}_{yu}(\omega)\hat{S}_u(\omega)^{-1},$$

where  $\hat{S}_{yu}, \hat{S}_u$  can be computed via DFT

# Periodogram analysis

Suppose an infinite-length discrete-time signal y(t) is windowed by a length-N window w(t),  $1 \le t \le N$ 

$$\tilde{y}(t) = w(t)y(t)$$

The Fourier transform of  $\tilde{y}(t)$  is then given by

$$Y_N(\omega) = \sum_{t=1}^N w(t)y(t)e^{-\mathrm{i}\omega t}$$

The *periodogram*, an estimate of  $S_y(\omega)$ , is obtained by

$$\hat{S}_y(\omega) = \frac{1}{CN} |Y_N(\omega)|^2,$$

where  $C = \frac{1}{N} \sum_{t=1}^{N} |w(t)|^2$  is a normalization factor

### Periodogram analysis

 $\hat{S}_y(\omega)$  is called *periodogram* when w(t) is rectangular, and *modified* periodogram for other types of windows, e.g., Hamming, Barlett, etc.

In practice, the periodogram is evaluated at a finite number of frequencies

$$\omega_k = 2\pi k/R, \quad 0 \le k \le R-1$$

by replacing  $\hat{S}_y(\omega)$  with the length-R DFT Y[k] of the length-N sequences y[k]:

$$\hat{S}_y(\omega_k) = \hat{S}_y[k] = \frac{1}{CN} |Y[k]|^2$$

- ullet Usually R>N to provide a finer resolution of the periodogram
- $C = (1/N) \sum_{t=1}^{N} |w(t)|^2$  is a normalization factor

Suppose we use a rectangular window of length N

$$\hat{S}_{y}(\omega) = \frac{1}{N} \sum_{n=1}^{N} \sum_{m=1}^{N} y(n) y^{*}(m) e^{-i\omega(m-n)}$$

$$= \frac{1}{N} \sum_{k=-N+1}^{N-1} \sum_{n=1}^{N-|k|} y(n+k) y^{*}(n) e^{-i\omega k}$$

$$= \sum_{k=-N+1}^{N-1} \hat{R}_{y}(k) e^{-i\omega k}$$

- ullet The periodogram is the Fourier transform of  $\hat{R}_y(k)$
- A few samples of y(n) is used in estimating  $\hat{R}_y(k)$  when k is large, yielding a poor estimate of  $R_y(k)$

Use the window functions that vanish for  $|\tau|>M$  to weight out the estimated correlation for large  $\tau$ 

Rectangular

$$w(\tau) = 1, \quad |\tau| \le M$$

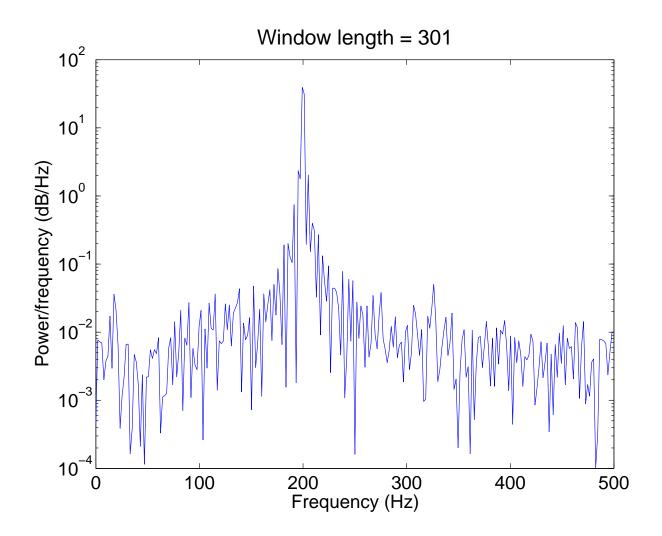
Barlett

$$w(\tau) = 1 - |\tau|/M, \quad |\tau| \le M$$

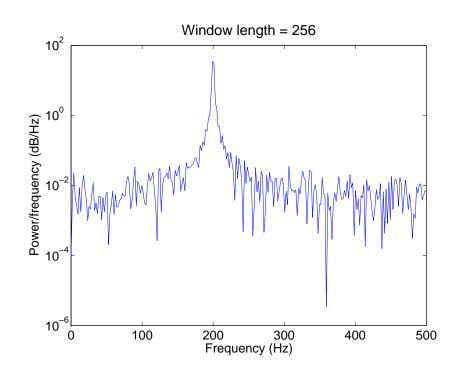
Hamming

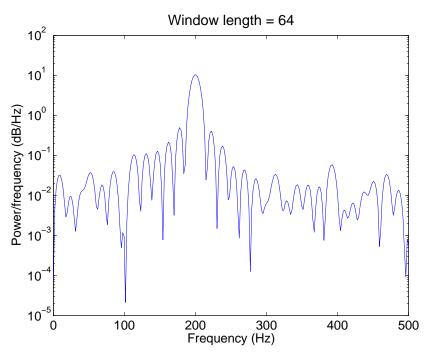
$$w(\tau) = 0.54 + 0.46 \cos\left(\frac{2\pi\tau}{2M+1}\right), \quad |\tau| \le M$$

 ${\cal M}$  should be small compared to  ${\cal N}$  to reduce the fluctuations of the periodogram



• 
$$y(t) = \cos(400\pi t) + \nu(t)$$
, with  $N = 301$ 





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